

3.9 Fluid dynamics

Ideal fluids^a

Continuity ^b	$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$	(3.285)	ρ density $\mathbf{v}$ fluid velocity field t time
Kelvin circulation	$\Gamma = \oint \mathbf{v} \cdot d\mathbf{l} = \text{constant}$	(3.286)	Γ circulation $d\mathbf{l}$ loop element
	$= \int_S \boldsymbol{\omega} \cdot d\mathbf{s}$	(3.287)	$d\mathbf{s}$ element of surface bounded by loop $\boldsymbol{\omega}$ vorticity ($= \nabla \times \mathbf{v}$)
Euler's equation ^c	$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{\nabla p}{\rho} + \mathbf{g}$	(3.288)	p pressure $\mathbf{g}$ gravitational field strength
	or $\frac{\partial}{\partial t}(\nabla \times \mathbf{v}) = \nabla \times [\mathbf{v} \times (\nabla \times \mathbf{v})]$	(3.289)	$(\mathbf{v} \cdot \nabla)$ advective operator
Bernoulli's equation (incompressible flow)	$\frac{1}{2} \rho v^2 + p + \rho g z = \text{constant}$	(3.290)	z altitude
Bernoulli's equation (compressible adiabatic flow) ^d	$\frac{1}{2} v^2 + \frac{\gamma}{\gamma - 1} \frac{p}{\rho} + g z = \text{constant}$	(3.291)	γ ratio of specific heat capacities (c_p/c_v)
	$= \frac{1}{2} v^2 + c_p T + g z$	(3.292)	c_p specific heat capacity at constant pressure T temperature
Hydrostatics	$\nabla p = \rho \mathbf{g}$	(3.293)	
Adiabatic lapse rate (ideal gas)	$\frac{dT}{dz} = -\frac{g}{c_p}$	(3.294)	

^aNo thermal conductivity or viscosity.

^bTrue generally.

^cThe second form of Euler's equation applies to incompressible flow only.

^dEquation (3.292) is true only for an ideal gas.

Potential flow^a

Velocity potential	$\mathbf{v} = \nabla \phi$	(3.295)	$\mathbf{v}$ velocity
	$\nabla^2 \phi = 0$	(3.296)	ϕ velocity potential
Vorticity condition	$\boldsymbol{\omega} = \nabla \times \mathbf{v} = 0$	(3.297)	$\boldsymbol{\omega}$ vorticity $\mathbf{F}$ drag force on moving sphere
Drag force on a sphere ^b	$\mathbf{F} = -\frac{2}{3} \pi \rho a^3 \dot{\mathbf{u}} = -\frac{1}{2} M_d \dot{\mathbf{u}}$	(3.298)	a sphere radius $\dot{\mathbf{u}}$ sphere acceleration ρ fluid density M_d displaced fluid mass

^aFor incompressible fluids.

^bThe effect of this drag force is to give the sphere an additional effective mass equal to half the mass of fluid displaced.

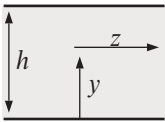
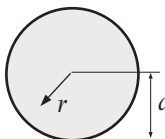
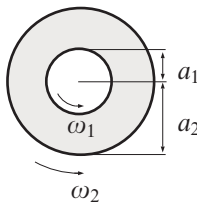
Viscous flow (incompressible)^a

Fluid stress	$\tau_{ij} = -p\delta_{ij} + \eta \left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right)$	(3.299)	τ_{ij} fluid stress tensor p hydrostatic pressure η shear viscosity v_i velocity along i axis δ_{ij} Kronecker delta
Navier–Stokes equation ^b	$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{\nabla p}{\rho} - \frac{\eta}{\rho} \nabla \times \boldsymbol{\omega} + \mathbf{g}$ $= -\frac{\nabla p}{\rho} + \frac{\eta}{\rho} \nabla^2 \mathbf{v} + \mathbf{g}$	(3.300) (3.301)	$\mathbf{v}$ fluid velocity field $\boldsymbol{\omega}$ vorticity $\mathbf{g}$ gravitational acceleration ρ density
Kinematic viscosity	$\nu = \eta / \rho$	(3.302)	ν kinematic viscosity

^aI.e., $\nabla \cdot \mathbf{v} = 0$, $\eta \neq 0$.

^bNeglecting bulk (second) viscosity.

Laminar viscous flow

Between parallel plates	$v_z(y) = \frac{1}{2\eta} y(h-y) \frac{\partial p}{\partial z}$	(3.303)	v_z flow velocity z direction of flow y distance from plate η shear viscosity p pressure	
Along a circular pipe ^a	$v_z(r) = \frac{1}{4\eta} (a^2 - r^2) \frac{\partial p}{\partial z}$ $Q = \frac{dV}{dt} = \frac{\pi a^4}{8\eta} \frac{\partial p}{\partial z}$	(3.304) (3.305)	r distance from pipe axis a pipe radius V volume	
Circulating between concentric rotating cylinders ^b	$G_z = \frac{4\pi\eta a_1^2 a_2^2}{a_2^2 - a_1^2} (\omega_2 - \omega_1)$	(3.306)	G_z axial couple between cylinders per unit length ω_i angular velocity of i th cylinder	
Along an annular pipe	$Q = \frac{\pi}{8\eta} \frac{\partial p}{\partial z} \left[a_2^4 - a_1^4 - \frac{(a_2^2 - a_1^2)^2}{\ln(a_2/a_1)} \right]$	(3.307)	a_1 inner radius a_2 outer radius Q volume discharge rate	

^aPoiseuille flow.

^bCouette flow.

Drag^a

On a sphere (Stokes's law)	$F = 6\pi a \eta v$	(3.308)	F drag force a radius
On a disk, broadside to flow	$F = 16a \eta v$	(3.309)	v velocity η shear viscosity
On a disk, edge on to flow	$F = 32a \eta v / 3$	(3.310)	

^aFor Reynolds numbers $\ll 1$.



Characteristic numbers

Reynolds number	$\text{Re} = \frac{\rho UL}{\eta} = \frac{\text{inertial force}}{\text{viscous force}} \quad (3.311)$	Re Reynolds number ρ density U characteristic velocity L characteristic scale-length η shear viscosity
Froude number ^a	$\text{F} = \frac{U^2}{Lg} = \frac{\text{inertial force}}{\text{gravitational force}} \quad (3.312)$	F Froude number g gravitational acceleration
Strouhal number ^b	$\text{S} = \frac{U\tau}{L} = \frac{\text{evolution scale}}{\text{physical scale}} \quad (3.313)$	S Strouhal number τ characteristic timescale
Prandtl number	$\text{P} = \frac{\eta c_p}{\lambda} = \frac{\text{momentum transport}}{\text{heat transport}} \quad (3.314)$	P Prandtl number c_p Specific heat capacity at constant pressure λ thermal conductivity
Mach number	$\text{M} = \frac{U}{c} = \frac{\text{speed}}{\text{sound speed}} \quad (3.315)$	M Mach number c sound speed
Rossby number	$\text{Ro} = \frac{U}{\Omega L} = \frac{\text{inertial force}}{\text{Coriolis force}} \quad (3.316)$	Ro Rossby number Ω angular velocity

^aSometimes the square root of this expression. L is usually the fluid depth.

^bSometimes the reciprocal of this expression.

Fluid waves

Sound waves	$v_p = \left(\frac{K}{\rho} \right)^{1/2} = \left(\frac{dp}{d\rho} \right)^{1/2} \quad (3.317)$	v_p wave (phase) speed K bulk modulus p pressure ρ density
In an ideal gas (adiabatic conditions) ^a	$v_p = \left(\frac{\gamma RT}{\mu} \right)^{1/2} = \left(\frac{\gamma p}{\rho} \right)^{1/2} \quad (3.318)$	γ ratio of heat capacities R molar gas constant T (absolute) temperature μ mean molecular mass
Gravity waves on a liquid surface ^b	$\omega^2 = gk \tanh kh \quad (3.319)$ $v_g \simeq \begin{cases} \frac{1}{2} \left(\frac{g}{k} \right)^{1/2} & (h \gg \lambda) \\ (gh)^{1/2} & (h \ll \lambda) \end{cases} \quad (3.320)$	v_g group speed of wave h liquid depth λ wavelength k wavenumber g gravitational acceleration ω angular frequency
Capillary waves (ripples) ^c	$\omega^2 = \frac{\sigma k^3}{\rho} \quad (3.321)$	σ surface tension
Capillary-gravity waves ($h \gg \lambda$)	$\omega^2 = gk + \frac{\sigma k^3}{\rho} \quad (3.322)$	

^aIf the waves are isothermal rather than adiabatic then $v_p = (p/\rho)^{1/2}$.

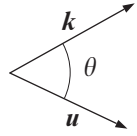
^bAmplitude $\ll$ wavelength.

^cIn the limit $k^2 \gg g\rho/\sigma$.



Doppler effect^a

Source at rest, observer moving at u	$\frac{v'}{v} = 1 - \frac{ u }{v_p} \cos \theta$	(3.323)	v', v'' observed frequency v emitted frequency v_p wave (phase) speed in fluid
Observer at rest, source moving at u	$\frac{v''}{v} = \frac{1}{1 - \frac{ u }{v_p} \cos \theta}$	(3.324)	u velocity θ angle between wavevector, k , and u



^aFor plane waves in a stationary fluid.

Wave speeds

Phase speed	$v_p = \frac{\omega}{k} = v\lambda$	(3.325)	v_p phase speed v frequency ω angular frequency ($= 2\pi v$) λ wavelength k wavenumber ($= 2\pi/\lambda$)
Group speed	$v_g = \frac{d\omega}{dk}$	(3.326)	v_g group speed
	$= v_p - \lambda \frac{dv_p}{d\lambda}$	(3.327)	

Shocks

Mach wedge ^a	$\sin \theta_w = \frac{v_p}{v_b}$	(3.328)	θ_w wedge semi-angle v_p wave (phase) speed v_b body speed
Kelvin wedge ^b	$\lambda_K = \frac{4\pi v_b^2}{3g}$	(3.329)	λ_K characteristic wavelength
	$\theta_w = \arcsin(1/3) = 19^\circ.5$	(3.330)	g gravitational acceleration
Spherical adiabatic shock ^c	$r \simeq \left(\frac{Et^2}{\rho_0} \right)^{1/5}$	(3.331)	r shock radius E energy release t time ρ_0 density of undisturbed medium
Rankine– Hugoniot shock relations ^d	$\frac{p_2}{p_1} = \frac{2\gamma M_1^2 - (\gamma - 1)}{\gamma + 1}$	(3.332)	1 upstream values 2 downstream values
	$\frac{v_1}{v_2} = \frac{\rho_2}{\rho_1} = \frac{\gamma + 1}{(\gamma - 1) + 2/M_1^2}$	(3.333)	p pressure v velocity
	$\frac{T_2}{T_1} = \frac{[2\gamma M_1^2 - (\gamma - 1)][2 + (\gamma - 1)M_1^2]}{(\gamma + 1)^2 M_1^2}$	(3.334)	T temperature ρ density γ ratio of specific heats M Mach number

^aApproximating the wake generated by supersonic motion of a body in a nondispersive medium.

^bFor gravity waves, e.g., in the wake of a boat. Note that the wedge semi-angle is independent of v_b .

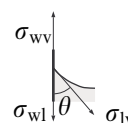
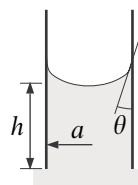
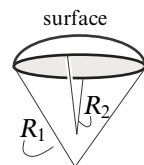
^cSedov–Taylor relation.

^dSolutions for a steady, normal shock, in the frame moving with the shock front. If $\gamma = 5/3$ then $v_1/v_2 \leq 4$.



Surface tension

Definition	$\sigma_{lv} = \frac{\text{surface energy}}{\text{area}}$	(3.335)	σ_{lv} surface tension (liquid/vapour interface)
	$= \frac{\text{surface tension}}{\text{length}}$	(3.336)	
Laplace's formula ^a	$\Delta p = \sigma_{lv} \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$	(3.337)	Δp pressure difference over surface R_i principal radii of curvature
Capillary constant	$c_c = \left(\frac{2\sigma_{lv}}{g\rho} \right)^{1/2}$	(3.338)	c_c capillary constant ρ liquid density g gravitational acceleration
Capillary rise (circular tube)	$h = \frac{2\sigma_{lv} \cos \theta}{\rho g a}$	(3.339)	h rise height θ contact angle a tube radius
Contact angle	$\cos \theta = \frac{\sigma_{wv} - \sigma_{wl}}{\sigma_{lv}}$	(3.340)	σ_{wv} wall/vapour surface tension σ_{wl} wall/liquid surface tension



^aFor a spherical bubble in a liquid $\Delta p = 2\sigma_{lv}/R$. For a soap bubble (two surfaces) $\Delta p = 4\sigma_{lv}/R$.

